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Search for $H \to \gamma\gamma$ produced in association with top quarks and constraints on the Yukawa coupling between the top quark and the Higgs boson using data taken at 7 TeV and 8 TeV with the ATLAS detector

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ABSTRACT

A search is performed for Higgs bosons produced in association with top quarks using the diphoton decay mode of the Higgs boson. Selection requirements are optimized separately for leptonic and fully hadronic final states from the top quark decays. The dataset used corresponds to an integrated luminosity of 4.5 fb^{-1} of proton–proton collisions at a center-of-mass energy of 7 TeV and 20.3 fb^{-1} at 8 TeV recorded by the ATLAS detector at the CERN Large Hadron Collider. No significant excess over the background prediction is observed and upper limits are set on the $t\bar{t}H$ production cross section. The observed exclusion upper limit at 95% confidence level is 6.7 times the predicted Standard Model cross section value. In addition, limits are set on the strength of the Yukawa coupling between the top quark and the Higgs boson, taking into account the dependence of the $t\bar{t}H$ and tH cross sections as well as the $H \to \gamma\gamma$ branching fraction on the Yukawa coupling. Lower and upper limits at 95% confidence level are set at −1.3 and +8.0 times the Yukawa coupling strength in the Standard Model.

1. Introduction

After the decades-long search for the Higgs boson [1–3], a particle consistent with the Standard Model (SM) Higgs boson has been discovered at the Large Hadron Collider (LHC) [4,5]. A notable property of the SM Higgs boson is its predicted large Yukawa coupling to top quarks, Y_t^{SM}. The measurement of Y_t is particularly important for understanding electroweak symmetry breaking and allows for testing theories beyond the SM (BSM).

The value of Y_t is indirectly tested by measurements sensitive to gluon fusion, ggF, the dominant Higgs boson production mechanism at the LHC, which receives large contributions from loop diagrams involving the top quark. In addition, Y_t is probed in the decay of the Higgs boson to two photons, $H \to \gamma\gamma$, as the decay width also involves loop diagrams with top quarks [6]. However, Y_t can be directly measured in the production of top–antitop quark pairs, $t\bar{t}$, in association with a Higgs boson [7–11], $t\bar{t}H$.

The production of the Higgs boson in association with a single top quark, tH,[1] is also sensitive to Y_t. Three processes contribute to tH production [12–16]: t-channel ($tHqb$) production, WtH production and s-channel tH production. The s-channel production is neglected in this Letter due to the much smaller cross section compared to $tHqb$ and WtH production. Examples of Feynman diagrams for $tHqb$ and WtH production are shown in Fig. 1.

In the SM, tH production is suppressed by the destructive interference between t-channel diagrams with Higgs bosons emitted from top quark and W boson lines, as for example shown in Fig. 1 (a) and Fig. 1 (b). In BSM theories [13–16], however, Y_t can have non-SM values, and in particular the relative sign between Y_t and g_{HWW}, which quantifies the coupling between the Higgs boson and the W boson, can be different from the SM prediction, which could lead to constructive instead of destructive interference in tH production. Hence, the tH production cross section is not only sensitive to the magnitude of Y_t but, in contrast to $t\bar{t}H$ production, it is also sensitive to the relative sign of Y_t with respect to g_{HWW}. A scale factor, κ_t, is introduced to describe the relation between Y_t and its SM value: $Y_t = \kappa_t Y_t^{\mathrm{SM}}$. Values of $\kappa_t \neq 1$ imply modifications of the Brout–Englert–Higgs mechanism and are assumed here to leave the top quark mass and decay properties unchanged. Furthermore, only SM particles are assumed to contribute to the decay width of the Higgs boson.

This Letter reports a search for $H \to \gamma\gamma$ in association with top quarks using data recorded with the ATLAS detector [18]. Measurements in the $H \to \gamma\gamma$ decay channel are challenging due to the

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[1] For simplicity, tH refers equally to $\bar{t}H$ in this Letter.

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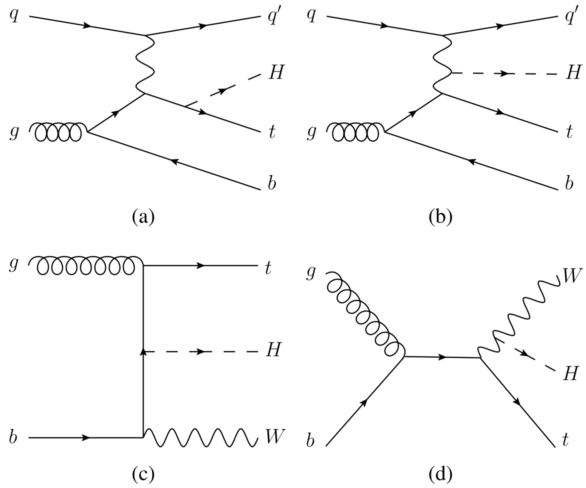

Fig. 1. Feynman diagrams showing examples for $tHqb$ (a, b) and WtH production (c, d). Higgs boson radiation off top quark and W boson lines is depicted. The $tHqb$ process is shown in the four-flavor scheme where no b-quarks are assumed to be present in the proton [17].

small branching fraction in the SM, $\mathrm{BR}(H \to \gamma\gamma) = 2.28 \times 10^{-3}$ for Higgs boson masses, m_H, around 125 GeV. However, the diphoton final state allows the diphoton invariant mass, $m_{\gamma\gamma}$, to be reconstructed with excellent resolution, strongly reducing the contribution from the backgrounds, which have a falling $m_{\gamma\gamma}$ spectrum, referred to as continuum background in the following. The contribution from the continuum background can be derived from data sidebands, thus not relying on theory assumptions. A previous search for $t\bar{t}H$ production by the CMS Collaboration has explored hadronic, diphoton and leptonic final states of the Higgs boson [19], setting an upper limit at the 95% confidence level (CL) on the ratio of the observed $t\bar{t}H$ production cross section to the SM expectation, called the signal strength $\mu_{t\bar{t}H}$, of 4.5.

This Letter also reports lower and upper limits at 95% CL on κ_t, taking into account the changes in the $t\bar{t}H$ and tH cross sections as well as the $H \to \gamma\gamma$ branching fraction [14–16]. BSM theories with values of $Y_t \neq Y_t^{\mathrm{SM}}$ are hence constrained.

2. The ATLAS detector

The ATLAS detector consists of an inner tracking detector system, electromagnetic and hadronic calorimeters, and an external muon spectrometer. Charged particles in the pseudorapidity[2] range $|\eta| < 2.5$ are reconstructed with the inner tracking detector, which is immersed in a 2 T axial field provided by a superconducting solenoid, and consists of pixel and microstrip semiconductor detectors, as well as a straw-tube transition radiation tracker. The solenoid is surrounded by sampling calorimeters, which span the pseudorapidity range up to $|\eta| = 4.9$. High-granularity liquid-argon (LAr) electromagnetic calorimeters are present up to $|\eta| = 3.2$. Hadronic calorimeters with scintillator tiles as active material cover $|\eta| < 1.74$, while LAr technology is used for hadronic calorimetry from $|\eta| = 1.5$ to $|\eta| = 4.9$. Outside the calorimeter system, air-core toroids provide a magnetic field for the muon spectrometer. Three stations of precision drift tubes and cathode strip chambers provide a measurements of muon tracks in the region $|\eta| < 2.7$. Resistive-plate and thin-gap chambers provide muon triggering capability up to $|\eta| < 2.4$. A detailed description of the ATLAS detector can be found in Ref. [18].

3. Data and Monte Carlo samples

3.1. Data samples

Data used for this analysis were recorded in pp collisions at $\sqrt{s} = 7$ TeV and 8 TeV in 2011 and 2012, respectively. All events satisfy data quality requirements ensuring proper functioning of the detector and trigger subsystems. The resulting datasets correspond to integrated luminosities of 4.5 fb^{-1} and 20.3 fb^{-1}, respectively [20]. For the 7 TeV dataset, events were triggered with a diphoton trigger with a threshold of 20 GeV on the transverse energy of each photon candidate. For the 8 TeV dataset, these thresholds were raised to 35 GeV for the highest-E_{T} (leading) photon candidate and 25 GeV for the second-highest-E_{T} (subleading) photon candidate.

3.2. Monte Carlo samples

The contribution from the continuum background is directly estimated from data. All processes involving $H \to \gamma\gamma$ decays, however, are estimated using Monte Carlo (MC) simulation samples.

The production of $t\bar{t}H$ events is modeled using next-to-leading-order (NLO) matrix elements obtained with the HELAC-One-loop package [21], where POWHEG-BOX [22–24] is interfaced to PYTHIA 8.1 [25] for showering and hadronization. CT10 [26] parton distribution functions (PDF) and the AU2 underlying event tune [27,28] are used. Production of $tHqb$ is simulated with MADGRAPH [29] in the four-flavor scheme with the CT10 PDF set, which provides a better description of the kinematics of the spectator b-quark than the five-flavor scheme [17]. PYTHIA 8.1 is used for showering and hadronization. Production of WtH is simulated in the five-flavor scheme by MADGRAPH5_AMC@NLO [30] interfaced to Herwig++ [31] using the CT10 PDF set. All tH samples are produced for three different values of κ_t: -1, 0 and $+1$. In the simulation of $t\bar{t}H$, $tHqb$ and WtH processes, diagrams with Higgs bosons radiated in the top quark decay are not taken into account because such contributions are negligible [32].

Higgs boson production by ggF and vector-boson fusion (VBF) is simulated with POWHEG-BOX [33,34] interfaced to PYTHIA 8.1 for showering and hadronization with CT10 PDF. Production of a Higgs boson in association with a W or Z boson (WH, ZH) is simulated with PYTHIA 8.1 using CTEQ6L1 [35] PDF.

All MC samples are generated at $m_H = 125$ GeV and are passed through a full GEANT4 [36] simulation of the ATLAS detector [37]. The simulated samples have additional pp collision events, pile-up, simulated by PYTHIA 8.1 added and weighted such that the average number of interactions per bunch-crossing is the same as in data.

The cross sections for $t\bar{t}H$ production were calculated at NLO in quantum chromodynamics (QCD) [7,9,38,39]. The cross sections for $tHqb$ production are calculated for different values of κ_t at LO using MADGRAPH with the renormalization and factorization scales set to 75 GeV, and with a minimum $p_{\mathrm{T},q}$ requirement of 10 GeV, consistent with the generated MC samples. LO-to-NLO K-factors are obtained by comparing the LO cross sections with the NLO cross sections calculated using MADGRAPH5_AMC@NLO. The cross sections for WtH production are calculated for different values of κ_t at NLO using MADGRAPH5_AMC@NLO with dynamic renormalization and factorization scales. Interference effects with $t\bar{t}H$ production are not considered, but are believed to be small given

[2] ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the centre of the detector and the z-axis along the beam pipe. The x-axis points from the IP to the centre of the LHC ring, and the y-axis points upward. Cylindrical coordinates (r, ϕ) are used in the transverse plane, ϕ being the azimuthal angle around the beam pipe. The pseudorapidity is defined in terms of the polar angle θ as $\eta = -\ln\tan(\theta/2)$. The transverse momentum is defined as $p_{\mathrm{T}} = p\sin\theta = p/\cosh\eta$, and the transverse energy E_{T} has an analogous definition.

Table 1
Production cross sections for the various Higgs boson processes at 7 TeV and 8 TeV before taking into account the BR($H \to \gamma\gamma$) at $m_H = 125$ GeV. Also quoted are the theoretical uncertainties from variations of the renormalization and factorization scales and uncertainties on the parton distribution functions [63,64].

Process	σ [pb] at 7 TeV	σ [pb] at 8 TeV
$t\bar{t}H$	$0.086^{+0.008}_{-0.011}$	$0.129^{+0.012}_{-0.016}$
$tHqb$, $\kappa_t = +1$	$0.0111^{+0.0009}_{-0.0008}$	$0.0172^{+0.0012}_{-0.0011}$
$tHqb$, $\kappa_t = 0$	$0.040^{+0.003}_{-0.003}$	$0.059^{+0.004}_{-0.004}$
$tHqb$, $\kappa_t = -1$	$0.129^{+0.010}_{-0.009}$	$0.197^{+0.014}_{-0.013}$
WtH, $\kappa_t = +1$	$0.0029^{+0.0007}_{-0.0006}$	$0.0047^{+0.0010}_{-0.0009}$
WtH, $\kappa_t = 0$	$0.0043^{+0.0011}_{-0.0008}$	$0.0073^{+0.0017}_{-0.0013}$
WtH, $\kappa_t = -1$	$0.016^{+0.004}_{-0.003}$	$0.027^{+0.006}_{-0.005}$
ggF	15.1 ± 1.6	19.3 ± 2.0
VBF	1.22 ± 0.03	1.58 ± 0.04
WH	0.579 ± 0.016	0.705 ± 0.018
ZH	0.335 ± 0.013	0.415 ± 0.017

that WtH is produced mostly without a second high-p_T b-quark in the final state.

The cross sections for ggF production were calculated at next-to-next-to leading order (NNLO) in QCD [40–45]. In addition, QCD soft-gluon resummation up to next-to-next-to-leading logarithms [46] is adopted to improve the NNLO calculation, and NLO electroweak (EW) corrections are applied [47,48]. The cross sections for VBF production were calculated including NLO QCD and EW corrections [49–51]. In addition, approximate NNLO QCD corrections are applied [52]. The cross sections for WH and ZH production were calculated at NLO [53] and NNLO [54] in QCD. Moreover, NLO EW corrections [55] are applied.

The theoretical uncertainties on the Higgs boson production cross sections come from varying the renormalization and factorization scales and from uncertainties on the parton distribution functions [26,56–58]. The Higgs boson decay branching fractions are taken from Refs. [59–62] and their uncertainties are compiled in Refs. [63,64]. A summary of the cross-section values and their uncertainties is given in Table 1.

4. Object and event selection

4.1. Object selection

Photons are reconstructed [65] from clusters of cells in the electromagnetic calorimeter in the region $|\eta| < 2.37$ excluding the transition region, $1.37 < |\eta| < 1.56$, between the barrel and endcap calorimeters. Unconverted photons are required to have no tracks associated with them; clusters from photons converted in the material between the production vertex and the calorimeter are allowed to have one or two associated tracks. The energies of the clusters are calibrated, separately for unconverted and converted photon candidates, in order to account for energy losses upstream of the calorimeter and for energy leakage outside of the cluster. Photons are required to pass a set of selection requirements on the reconstructed shower shape as well as the following isolation requirements: the sum of the p_T of all particles featuring tracks with $p_T > 1$ GeV in a cone of size $\Delta R \equiv \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2} = 0.2$ around the photon is required to be smaller than 2.6 (2.2) GeV for the $\sqrt{s} = 8$ TeV (7 TeV) data. Tracks from converted photons are excluded from the sum. Moreover, the sum of the E_T values in the calorimeter cells in a cone of size $\Delta R = 0.4$ around the photon is required to be smaller than 6 (5.5) GeV for the 8 TeV (7 TeV) data. The calorimeter isolation is corrected for photon energy leakage. It is also corrected event-by-event by using the ambient energy from pile-up and the underlying event [66,67]. Only events with two photons are retained and a diphoton vertex is reconstructed by a neural-network-based algorithm [68], which uses as input the trajectories of the two photons and the tracks associated with different vertex candidates. The photon trajectory is determined from the longitudinal profile of the photon shower in the calorimeter, the average pp collision point, and for converted photons from the direction of the associated tracks. The leading (subleading) photon is required to have $E_T > 0.35 \times m_{\gamma\gamma}$ ($0.25 \times m_{\gamma\gamma}$), and the diphoton mass is required to be between 105 GeV and 160 GeV.

Electrons are reconstructed [69] from clusters of cells in the electromagnetic calorimeter with an associated track. Only clusters in the region $|\eta| < 2.47$ are considered and are required to fulfill requirements on their shape to be consistent with an electron. The electron E_T has to be larger than 15 GeV. In addition, electrons must be isolated: the E_T in a cone of size $\Delta R = 0.4$ around the electron and the sum of the transverse momenta of the tracks in a cone of size $\Delta R = 0.2$ around the electron must be smaller than 20% and 15% of the electron E_T, respectively.

Muons are reconstructed [70] by combining tracks in the inner detector with tracks or track-segments in the muon spectrometer. Muons are required to satisfy $|\eta| < 2.7$ and $p_T > 10$ GeV and have to be isolated: muons closer than $\Delta R = 0.4$ to a jet or to one of the two photons are not considered. Moreover, the E_T in a cone of size $\Delta R = 0.4$ around the muon and the sum of the transverse momenta of the tracks in a cone of size $\Delta R = 0.2$ around the muon must be smaller than 20% and 15% of the muon p_T, respectively.

Jets are reconstructed from clusters of cells in the calorimeter with the anti-k_t algorithm [71] with a radius parameter of 0.4. They are calibrated to the hadronic energy scale [72], and only those with $p_T > 25$ GeV and $|\eta| < 2.5$ are considered. The jet energy is corrected for energy deposits from additional soft interactions in the event [73]. In order to suppress jets from additional interactions, the jet vertex fraction (JVF) must be larger than 50% for jets with $p_T < 50$ GeV and $|\eta| < 2.4$. The JVF is defined from the summed track p_T as the fraction associated with the primary diphoton vertex, where all tracks with $p_T > 0.5$ GeV matched to the jet are considered.

Jets containing b-quarks are identified with a neural-network-based b-tagging algorithm, which combines variables from impact parameter, secondary vertex and decay topology algorithms evaluating the track parameters associated with the jet [74]. Three different working points (WP) with efficiencies of 60%, 70% and 80% for identifying b-jets are used for 8 TeV data. For 7 TeV data, a slightly different optimization of the b-tagging algorithm with a WP corresponding to an efficiency of 85% is used. The b-tagging and mistagging efficiencies are measured in data using dijet and $t\bar{t}$ events [75].

The magnitude of the missing transverse momentum in each event, E_T^{miss}, is calculated using clusters of cells in the calorimeter. Corrections are applied for identified photons, electrons, muons and jets according to special E_T^{miss} object identification requirements [76].

In order to avoid double-counting of reconstructed objects, electrons with a distance in η–ϕ space smaller than 0.4 to one of the two photons, $\Delta R(e,\gamma)$, are not considered. In addition, jets with $\Delta R(\text{jet},\gamma) < 0.4$ or $\Delta R(\text{jet}, e) < 0.2$ are removed.

4.2. Event selection

In addition to the requirement of two good photons satisfying the criteria described in Section 4.1, two different event selections were optimized in order to efficiently select leptonic $t\bar{t}H$ events (leptonic category) as well as all-hadronic $t\bar{t}H$ events (hadronic category). The optimization targeted an optimal expected limit on the signal strength $\mu_{t\bar{t}H}$ in case no evidence for $t\bar{t}H$ production is found. However, the requirements for the leptonic category are

Table 2
Expected numbers of $H \to \gamma\gamma$ events (N_H) from an SM Higgs boson with $m_H = 125.4$ GeV after the event selection. These combined yields are normalized to 4.5 fb^{-1} for the 7 TeV data and to 20.3 fb^{-1} for the 8 TeV data, and are listed in the table along with the percent contribution of each Higgs boson production process with respect to the sum of all Higgs boson production processes. The numbers of fitted continuum background events (N_B) for the 7 TeV and 8 TeV data are also shown, where N_B is the integral of the continuum background in the $m_{\gamma\gamma}$ range 120–130 GeV, which is determined by an unbinned signal-plus-background fit to all categories with one common scale factor for the $H \to \gamma\gamma$ normalization. The uncertainty on N_B is the statistical uncertainty calculated from $\delta N_B = \delta N_{\text{tot}} N_B / N_{\text{tot}}$, where N_{tot} is the total number of background events in the full $m_{\gamma\gamma}$ range 105–160 GeV estimated from an unbinned signal-plus-background likelihood fit, and δN denotes the Poisson uncertainty on N.

Category	N_H	ggF	VBF	WH	ZH	$t\bar{t}H$	$tHqb$	WtH	N_B
7 TeV leptonic selection	0.10	0.6	0.1	14.9	4.0	72.6	5.3	2.5	$0.5^{+0.5}_{-0.3}$
7 TeV hadronic selection	0.07	10.5	1.3	1.3	1.4	80.9	2.6	1.9	$0.5^{+0.5}_{-0.3}$
8 TeV leptonic selection	0.58	1.0	0.2	8.1	2.3	80.3	5.6	2.6	$0.9^{+0.6}_{-0.4}$
8 TeV hadronic selection	0.49	7.3	1.0	0.7	1.3	84.2	3.4	2.1	$2.7^{+0.9}_{-0.7}$

kept loose enough in order to also allow high selection efficiency for $tHqb$ and WtH production.

In this analysis, we assume that the top quark only decays to a W boson and a b-quark. The leptonic selection targets both the single-lepton decays of the $t\bar{t}$ pairs, where one of the W bosons decays leptonically and the other one decays hadronically, and the dilepton decays of $t\bar{t}$ pairs, where both W bosons decay leptonically. Events are selected by requiring at least one electron or muon, at least one b-tagged jet using the 80% (85%) WP for 8 TeV (7 TeV) data and $E_{\text{T}}^{\text{miss}} > 20$ GeV. The $E_{\text{T}}^{\text{miss}}$ requirement is imposed to reduce backgrounds from final states without top quarks and it is not used for events with two or more b-tagged jets. Events with an electron–photon invariant mass in the range 84–94 GeV are rejected in order to reduce the background contribution from $Z \to ee$ events with one electron misidentified as a photon.

The hadronic selection targets events where both W bosons, from the top quark decays, decay hadronically. No electrons or muons may be identified in the event. Events must fulfill requirements on the number of jets and the number of b-tagged jets. For the 8 TeV dataset three sets of requirements are defined, out of which at least one must be satisfied for an event to be considered:

1. At least six jets, out of which at least two must be b-tagged using the 80% WP.
2. At least five jets with an increased p_{T} threshold of 30 GeV, out of which at least two must be b-tagged using the 70% WP.
3. At least six jets with an increased p_{T} threshold of 30 GeV, out of which at least one must be b-tagged using the 60% WP.

These requirements were optimized to suppress in particular the contribution from ggF Higgs boson production with $H \to \gamma\gamma$ to the hadronic category, while retaining good sensitivity to $t\bar{t}H$ production. For the 7 TeV dataset only events with at least six jets, at least two of which are b-tagged with the 85% WP, are considered.

Table 2 summarizes the expected numbers of events in each category for $m_H = 125.4$ GeV, the Higgs boson mass measured by the ATLAS Collaboration [68]. The breakdown into the different Higgs boson production processes is given. The combined selection efficiencies in the 7 TeV and 8 TeV data for $t\bar{t}H$ production at $m_H = 125.4$ GeV are approximately 14.6% and 14.8%, respectively. For SM $tHqb$ (WtH) production the combined selection efficiencies for 7 TeV and 8 TeV are approximately 6.2% (12.9%) and 6.2% (11.9%), respectively.

5. Analysis

In order to separate processes involving $H \to \gamma\gamma$ decays from the continuum background, a localized excess of events is searched for in the $m_{\gamma\gamma}$ spectrum around $m_H = 125.4$ GeV. Probability distribution functions for the $H \to \gamma\gamma$ resonance and continuum background $m_{\gamma\gamma}$ distributions are defined in the range of 105–160 GeV as described below, and the numbers of Higgs boson and continuum background events are estimated from an unbinned signal-plus-background likelihood fit to the full $m_{\gamma\gamma}$ distributions in the leptonic and hadronic categories. Systematic uncertainties are taken into account as nuisance parameters, which are fitted within their external constraints.

The sum of a Crystal Ball function [77] and a Gaussian function is used to describe the $m_{\gamma\gamma}$ distribution from $H \to \gamma\gamma$ decays obtained from MC simulations [78]. The Gaussian function accounts only for a small fraction of the total $H \to \gamma\gamma$ resonance signal, describing small tails of the shape which cannot be characterized by the Crystal Ball function. The parameters of these functions are interpolated between the values fitted to a series of MC samples generated in steps of 5 GeV in m_H, in order to allow for the evaluation of the resonance shape for intermediate masses including $m_H = 125.4$ GeV, where MC samples are not available. The relative fraction of the Gaussian component with respect to the full $H \to \gamma\gamma$ resonance shape is not varied as a function of m_H. Shapes with different parameter values are defined for the 7 TeV and 8 TeV data. The $m_{\gamma\gamma}$ resolution, which is quantified by half of the smallest $m_{\gamma\gamma}$ interval containing 68% of the signal events, is 1.42 GeV for the 7 TeV data and 1.56 GeV for the 8 TeV data in the leptonic categories. The values in the hadronic categories are consistent with the ones in the leptonic categories within statistical uncertainties. The small difference in $m_{\gamma\gamma}$ resolution between 7 TeV and 8 TeV is due to a difference in the effective constant term for the calorimeter energy resolution and due to the lower level of pile-up in the 7 TeV data [68]. The $m_{\gamma\gamma}$ resolution is dominated by the photon energy resolution. The small change in acceptance for $t\bar{t}H$ production is interpolated using MC samples generated with different hypothesized values of m_H also. For all other Higgs boson production processes, the difference in acceptance between $m_H = 125$ GeV and $m_H = 125.4$ GeV is found to be negligible.

An exponential function, $e^{a\, m_{\gamma\gamma}}$, with $a \le 0$ is chosen for both categories as a model for the continuum background following the method previously used in Ref. [5]. The choice of fit function is validated in data control regions obtained by loosening the photon identification and isolation requirements. These control regions are dominated by jets misidentified as photons, and the systematic uncertainties derived from these control regions (cf. Section 6) are hence only approximate. In both the leptonic and the hadronic category, the same continuum background shape is used for 7 TeV and 8 TeV data, because the 7 TeV data alone is not expected to strongly constrain the parameter a given the expected low number of events.

In the range 105 GeV $< m_{\gamma\gamma} <$ 160 GeV, 3 (3) events are found in the leptonic (hadronic) category in the 7 TeV and 5 (15) events are found in the 8 TeV data. The results of the fits for the leptonic and hadronic categories are shown in Fig. 2, separately for 7 TeV and 8 TeV data. The fitted numbers of continuum background events in a window of 120–130 GeV are shown in Table 2.

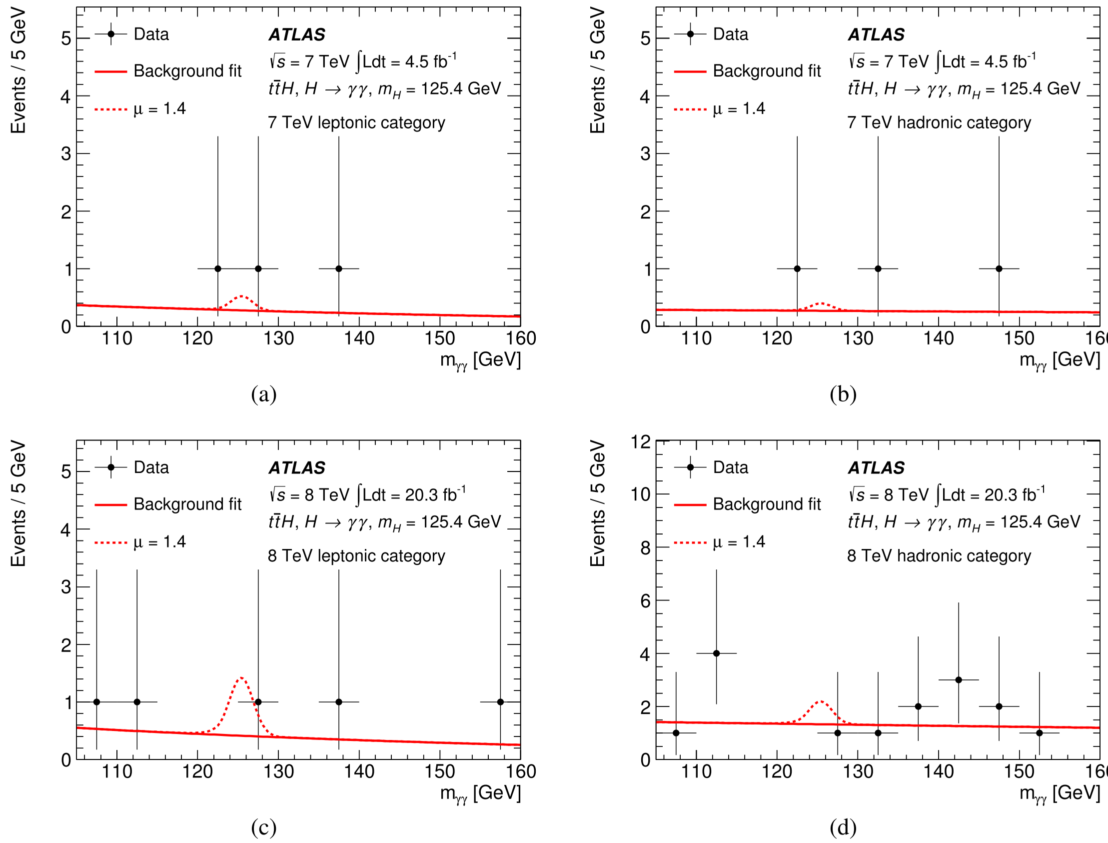

Fig. 2. Distributions of the diphoton invariant mass, $m_{\gamma\gamma}$, for the leptonic (left) and hadronic (right) category for data at 7 TeV (top) and data at 8 TeV (bottom). An unbinned signal-plus-background likelihood fit to the full spectra is used to estimate the number of events from continuum background (solid line) as well as from SM Higgs boson production (dashed line). The signal strength, μ, is a parameter common to all categories and its best-fit value is $\mu = 1.4$ for $m_H = 125.4$ GeV.

Table 3
Summary of systematic uncertainties on the final yield of events for 8 TeV data from $t\bar{t}H$, $tHqb$ and WtH production after applying the leptonic and hadronic selection requirements. The uncertainties are also shown for other Higgs boson production processes that do not include the associated production of top quarks and have significant contributions to the event selection. These are WH production in the leptonic category and ggF production in the hadronic category. For both tH production processes, the maximum uncertainty observed for all values of κ_t generated (+1, 0, −1) is reported.

	$t\bar{t}H$ [%]		$tHqb$ [%]		WtH [%]		ggF [%]	WH [%]
	had.	lep.	had.	lep.	had.	lep.	had.	lep.
Luminosity	±2.8							
Photons	±5.6	±5.5	±5.6	±5.5	±5.6	±5.5	±5.6	±5.5
Leptons	< 0.1	±0.7	< 0.1	±0.6	< 0.1	±0.6	< 0.1	±0.7
Jets and E_T^{miss}	±7.4	±0.7	±16	±1.9	±11	±2.1	±29	±10
Bkg. modeling	0.24 evt.	0.16 evt.	applied on the sum of all Higgs boson production processes					
Theory ($\sigma \times$ BR)	+10, −13		+7, −6		+14, −12		+11, −11	+5.5, −5.4
MC modeling	±11	±3.3	±12	±4.4	±12	±4.6	±130	±100

6. Systematic uncertainties

Systematic uncertainties from various sources affect both the expected number of events for different Higgs boson production processes and the $m_{\gamma\gamma}$ resonance shape. An overview of all systematic uncertainties for 8 TeV data is shown in Table 3 for $t\bar{t}H$, $tHqb$ and WtH production. The uncertainties are also shown for other Higgs boson production processes that do not include the associated production of top quarks and have significant contributions to the event selection. These are WH production in the leptonic category and ggF production in the hadronic category.

The uncertainty on the integrated luminosity is 2.8% (1.8%) for the 8 TeV (7 TeV) data as derived following the same methodology as that detailed in Ref. [20] using beam-separation scans. For 8 TeV data, the trigger efficiency [79] was measured to be $99.5 \pm 0.2\%$. For 7 TeV data, the efficiency was measured to be compatible with 100% within an uncertainty of 0.2%. The uncertainty in the combined diphoton identification efficiency is 1.0% (8.4%) [80] for 8 TeV (7 TeV) data. Due to the high jet multiplicity in this analysis an additional uncertainty of 4% is added to account for possible mismodeling of the photon identification efficiency. This additional uncertainty is obtained from data–MC comparisons of electron efficiencies in $Z(\to ee)$ + jets events, where photon identification requirements are applied to the electron clusters [81]. Analogously, an additional uncertainty of 3% is assessed for the efficiency of the combined diphoton isolation requirement, and is added in quadrature to the nominal uncertainty of 2.3% (2.1%) in the hadronic (leptonic) category. The uncertainty on the photon energy scale [80] was found to have a negligible effect on the expected yields. Its effect on the peak position, however, is taken into account, but has a negligible impact on the results. The uncertainty in the photon energy resolution translates into an uncertainty on the $m_{\gamma\gamma}$

resolution, and is based on the resolution measured with $Z \to ee$ events [80]. The total $m_{\gamma\gamma}$ resolution uncertainty is 12% for both the 7 TeV and 8 TeV dataset, which is less than 0.2 GeV.

The uncertainties due to the lepton reconstruction, identification, isolation, and energy/momentum scale and resolution combine to less than 1% for all channels. Uncertainties on the jet energy scale are taken into account, as well as uncertainties on the jet energy resolution, and on the modeling of the JVF and of the b-tagging efficiencies. All object uncertainties which change the energy or momentum of the corresponding objects are propagated to the $E_{\mathrm{T}}^{\mathrm{miss}}$ calculation, and additional uncertainties are taken into account for energy deposits which only enter the $E_{\mathrm{T}}^{\mathrm{miss}}$ calculation, but are not part of other objects.

Systematic uncertainties due to the choice of the continuum background fit model are estimated by fitting continuum background distributions in control regions with a Higgs boson plus continuum background model and quantifying the apparent number of Higgs boson events introduced [5]. The systematic uncertainty is chosen to be the maximal apparent number of Higgs boson events in a narrow mass range around 125.4 GeV. Since the contributions from different background processes in the control region may be different from their contributions in the four categories, the estimate of this uncertainty is approximate, but its impact on the final results is very small. An uncertainty of 0.24 (0.16) events is estimated in the 8 TeV hadronic (leptonic) category as the apparent number of Higgs boson events under the Higgs boson peak. For the 7 TeV dataset, uncertainties of 0.12 and 0.01 events are estimated, where all of these numbers have a non-negligible statistical component from the limited number of events in the control regions considered. The number of events is lowest in the control region for the hadronic category in 7 TeV data (266 events).

The theoretical uncertainties on the different Higgs boson production cross sections due to uncertainties in the PDF, missing higher-order perturbative QCD corrections estimated by varying the renormalization and factorization scales, and the BR($H \to \gamma\gamma$) are detailed in Refs. [26,56–58,62–64,82].

Additional uncertainties are included in "MC modeling" in Table 3. These take into account changes in the acceptance when the renormalization and factorization scales are varied, an uncertainty on the modeling of the underlying event, which is conservatively estimated by comparing MC samples with and without multiple parton scattering, and an uncertainty due to the limited number of events present in the MC samples after the event selection and categorization are applied. Moreover, uncertainties of 100% are assigned to the expected numbers of events from ggF, VBF and WH production in association with additional b-jets. The size of these uncertainties is motivated by recent measurements of $t\bar{t}$ and vector-boson production in association with b-jets [83,84].

7. Results

In total, 5 candidate events with $m_{\gamma\gamma}$ in the range 120–130 GeV are found in the leptonic and hadronic categories. The total expected yield of Higgs boson production is 1.3 events compared to a continuum background of $4.6^{+1.3}_{-0.9}$ events (see Table 2). The $m_{\gamma\gamma}$ spectra for the candidate events are shown in Fig. 2 together with the fitted continuum background and the total contribution from $H \to \gamma\gamma$ processes, where the signal strength, μ, is a parameter common to all four categories. The best-fit signal strength for all $H \to \gamma\gamma$ processes together is $1.4^{+2.1}_{-1.4}$(stat.)$^{+0.6}_{-0.3}$(syst.), where the quoted overall systematic uncertainty is derived by quadratically subtracting the statistical uncertainty from the total uncertainty. When the yields for all $H \to \gamma\gamma$ processes, including tH production but not $t\bar{t}H$ production, are set to their respective SM ex-

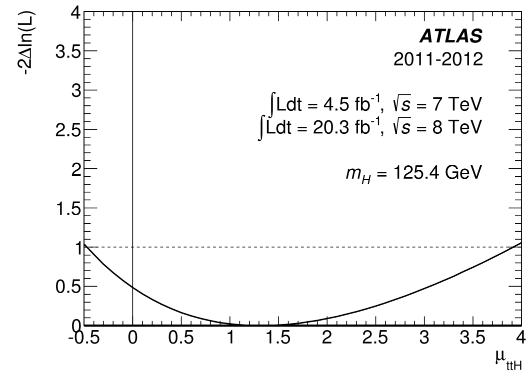

Fig. 3. Negative log-likelihood scan for the $t\bar{t}H$ cross section times BR($H \to \gamma\gamma$) relative to the SM expectation, $\mu_{t\bar{t}H}$, at $m_H = 125.4$ GeV, where all other Higgs boson production cross sections, including the cross section for tH production, are set to their respective SM expectations.

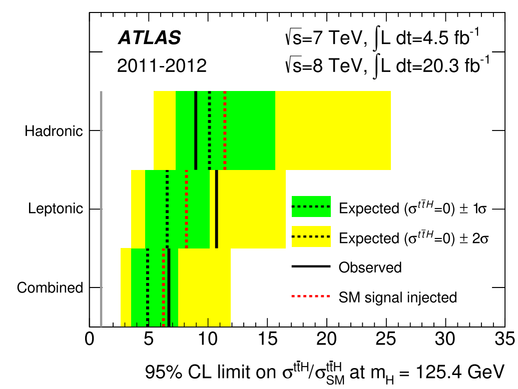

Fig. 4. Observed and expected 95% CL upper limits on the $t\bar{t}H$ production cross section times BR($H \to \gamma\gamma$). All other Higgs boson production cross sections, including the cross section for tH production, are set to their respective SM expectations. While the expected limits are calculated for the case where $t\bar{t}H$ production is not present, the lines denoted by "SM signal injected" show the expected 95% CL limits for a dataset corresponding to continuum background plus SM Higgs boson production. The limits are given relative to the SM expectations and at $m_H = 125.4$ GeV.

pected number of events, a best-fit value of $1.3^{+2.5}_{-1.7}$(stat.)$^{+0.8}_{-0.4}$(syst.) is obtained for $\mu_{t\bar{t}H}$, which is also shown in the scan of the likelihood in Fig. 3. This best-fit value of $\mu_{t\bar{t}H}$ is consistent with the SM expectation of one, but does not represent a significant excess over the predicted background rate, and CL_s-based [85] 95% CL exclusion upper limits are set for $t\bar{t}H$ production times BR($H \to \gamma\gamma$). Limits are set using the asymptotic formulae discussed in Ref. [86] with the profile likelihood ratio as test statistic. The results are found to be consistent with limits derived from ensembles of pseudo-experiments. The observed and expected upper limits for $\mu_{t\bar{t}H}$ at $m_H = 125.4$ GeV are summarized in Fig. 4 as well as in Table 4, where the expected limits assume $\mu_{t\bar{t}H} = 0$. The non-$t\bar{t}H$ Higgs boson production modes, including tH, are fixed to their SM expectations with corresponding theory and experimental uncertainties assigned. An upper limit of 6.7 times the SM cross section times BR($H \to \gamma\gamma$) is observed. Upper limits at 95% CL are also set on the signal strength of the sum of all $H \to \gamma\gamma$ processes, μ, and the observed (expected) limit is 5.7 (3.8).

These results are also interpreted as 95% CL limits on the strength parameter κ_t of the top quark–Higgs boson Yukawa coupling. Variations in κ_t not only change the production cross sections of the $t\bar{t}H$ and tH processes, but also affect BR($H \to \gamma\gamma$), and the cross sections of the other Higgs boson production processes [82]. Fig. 5 illustrates the dependence of the $t\bar{t}H$ and tH cross sections and of the BR($H \to \gamma\gamma$) on κ_t. For $\kappa_t = 0$, the $t\bar{t}H$

Table 4
Observed and expected 95% CL upper limits on the $t\bar{t}H$ production cross section times BR($H \to \gamma\gamma$) relative to the SM cross section times BR($H \to \gamma\gamma$) at $m_H = 125.4$ GeV. All other Higgs boson production cross sections, including the cross section for tH production, are set to their respective SM expectations. In addition, the expected limits corresponding to $+2\sigma$, $+1\sigma$, -1σ, and -2σ variations are shown. The expected limits are calculated for the case where $t\bar{t}H$ production is not present. The results are given for the combination of leptonic and hadronic categories with all systematic uncertainties included, and also for leptonic and hadronic categories separately, as well as for the expected limits additionally with only statistical uncertainties considered.

	Observed limit	Expected limit	$+2\sigma$	$+1\sigma$	-1σ	-2σ
Combined (with systematics)	6.7	4.9	11.9	7.5	3.5	2.6
Combined (statistics only)		4.7	10.5	7.0	3.4	2.5
Leptonic (with systematics)	10.7	6.6	16.5	10.1	4.7	3.5
Leptonic (statistics only)		6.4	15.1	9.6	4.6	3.4
Hadronic (with systematics)	9.0	10.1	25.4	15.6	7.3	5.4
Hadronic (statistics only)		9.5	21.4	14.1	6.8	5.1

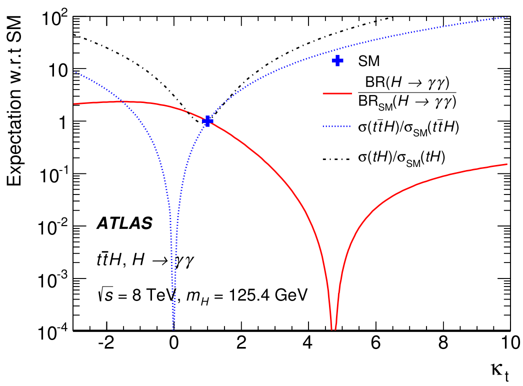

Fig. 5. Production cross sections for $t\bar{t}H$ and tH divided by their SM expectations as a function of the scale factor to the top quark–Higgs boson Yukawa coupling, κ_t. Production of tH comprises the $tHqb$ and WtH processes. Also shown is the dependence of the BR($H \to \gamma\gamma$) with respect to its SM expectation on κ_t.

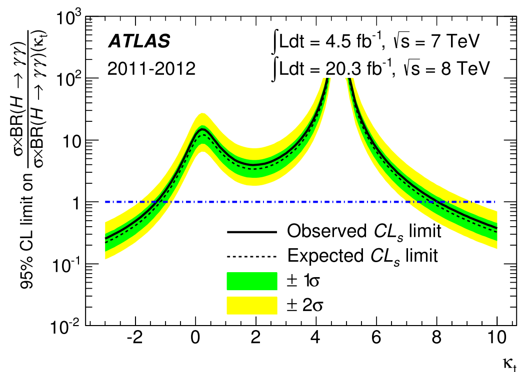

Fig. 6. Observed and expected 95% CL upper limits on the inclusive Higgs boson production cross section with respect to the cross section times BR($H \to \gamma\gamma$) for different values of κ_t at $m_H = 125.4$ GeV, where κ_t is the strength parameter for the top quark–Higgs boson Yukawa coupling. All Higgs boson production processes are considered for the inclusive production cross section. The expected limits are calculated for the case where $\kappa_t = +1$. The CL_s alternative hypothesis is given by continuum background plus SM Higgs boson production.

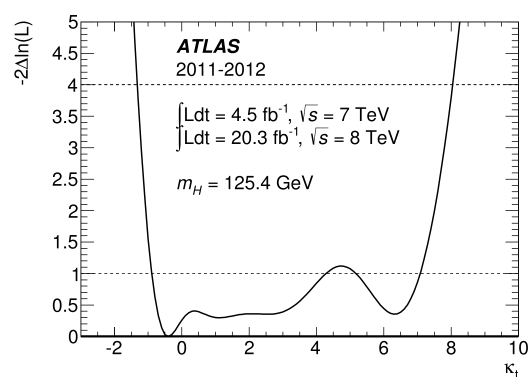

Fig. 7. Negative log-likelihood scan of κ_t at $m_H = 125.4$ GeV, where κ_t is the strength parameter for the top quark–Higgs boson Yukawa coupling.

process is turned off, and the top quark contribution to tH production and to the loop-induced $H \to \gamma\gamma$ decay is removed, leaving mainly the contribution from W bosons. For values of $\kappa_t < 0$, on the other hand, the interference between contributions from W bosons and top quarks to tH production and to the BR($H \to \gamma\gamma$) becomes constructive, thus enhancing the two processes with respect to their respective SM expectations. Cancellations of the contributions of top quarks and W bosons to the loop-induced $H \to \gamma\gamma$ decay lead to a minimum of the BR($H \to \gamma\gamma$) around a value of $\kappa_t = +4.7$. The combined selection efficiency differs slightly for the three values of κ_t for which $tHqb$ and WtH MC samples were generated. From these, the efficiency at different values of κ_t in the range $[-3, +10]$ is calculated by combining reweighted MC samples with $\kappa_t = +1, 0$ and -1. The weight for each sample is assigned in such a way that the cross-section value from the combination follows the prediction shown in Fig. 5. The largest relative difference with respect to the efficiency at $\kappa_t = +1$ over the entire range is found to be 14% (20%) for $tHqb$ (WtH) production.

All $H \to \gamma\gamma$ processes are considered and 95% CL limits are set on the total Higgs boson production cross section times BR($H \to \gamma\gamma$) with respect to the SM cross section for different values of κ_t. Coupling strengths other than κ_t are set to their respective SM values. The continuum background plus SM Higgs boson production ($\kappa_t = +1$) is taken as alternative hypothesis.

The observed and expected limits on κ_t at $m_H = 125.4$ GeV are summarized in Fig. 6, where the observed (expected) lower and upper limits on κ_t at 95% CL are -1.3 and $+8.0$ (-1.2 and $+7.8$). The expected limits assume $\kappa_t = +1$. The form of the limit curve shown in Fig. 6 is the result of the different dependencies of the different Higgs boson production processes as well as the BR($H \to \gamma\gamma$) on κ_t. The negative log-likelihood scan of κ_t is shown in Fig. 7 and it shows that the data are consistent with the SM expectation of $\kappa_t = +1$. Although two different values of κ_t exist with the same total number of expected events, there are no double minima at zero shown in Fig. 6 because different relative contributions from the Higgs boson production processes in different categories have lifted the degeneracy of the likelihood.

8. Conclusion

A search for Higgs boson production in association with top quarks in the $H \to \gamma\gamma$ decay channel is presented using leptonic and hadronic $t\bar{t}$ decays. Data at 7 TeV and 8 TeV corresponding to 4.5 fb^{-1} and 20.3 fb^{-1} taken in pp collisions with the ATLAS detector at the LHC were analyzed. No significant excess over the background prediction is observed and upper limits at 95% CL are set on the $t\bar{t}H$ production cross section. The observed exclusion limit at $m_H = 125.4$ GeV is found to be 6.7 times the predicted SM cross section. The corresponding lower and upper limits on the top quark–Higgs boson Yukawa coupling strength parameter κ_t are found to be -1.3 and $+8.0$, which in particular constrain models with a negative sign of the coupling.

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